

# Aharonov-Bohm Effect for pedestrian

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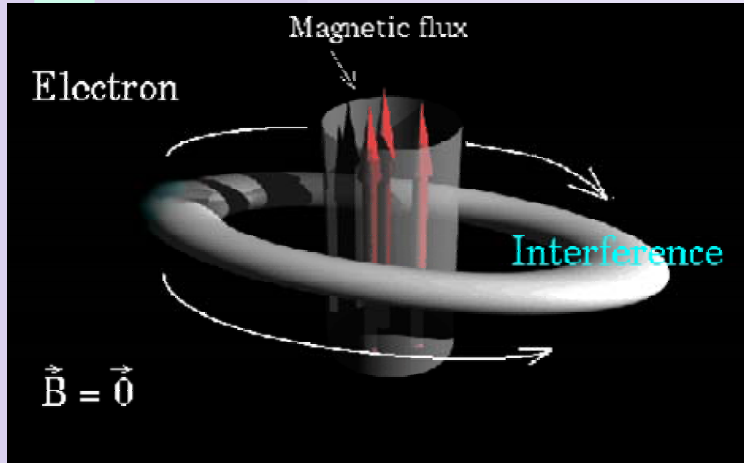
Castelló (SPAIN), June 2005



**Commemorative Meeting in honour of Professor Brian G. Wybourne**

# Aharonov-Bohm Effect

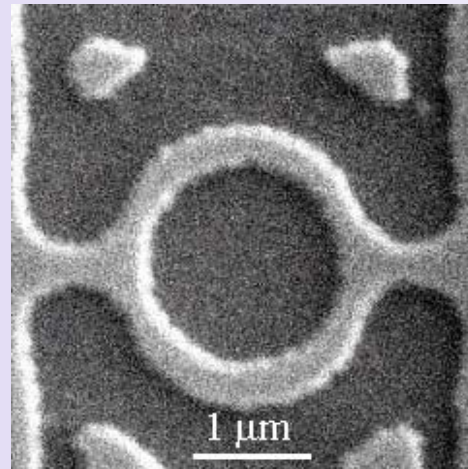
Y. Aharonov, D. Bohm, *Phys. Rev.* **115** (1959) 485



- Classical mechanics: equations of motion can always be expressed in term of field alone.
- Quantum mechanics: canonical formalism. Potentials cannot be eliminated.
- An electron can be influenced by the potentials even if no fields acts upon it.
- $\oint Adl$  Gauge independent

# Semiconductor Quantum Rings

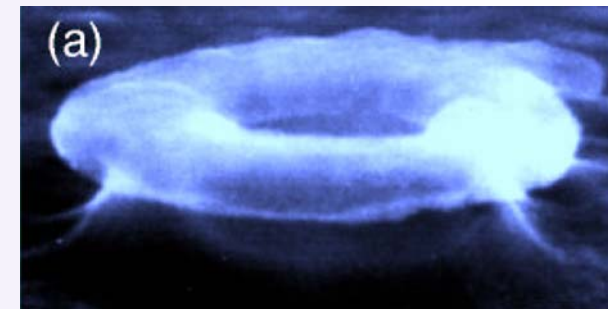
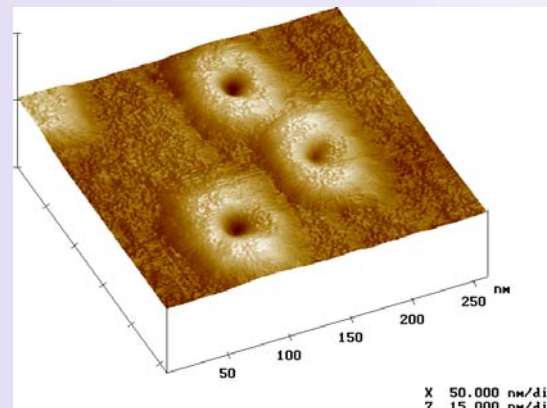
Litographic rings  
GaAs/AlGaAs



*A. Fuhrer et al., Nature*  
**413** (2001) 822;

*M. Bayer et al.,*  
*Phys. Rev. Lett.*  
**90** (2003) 186801.

self-assembled rings  
InAs



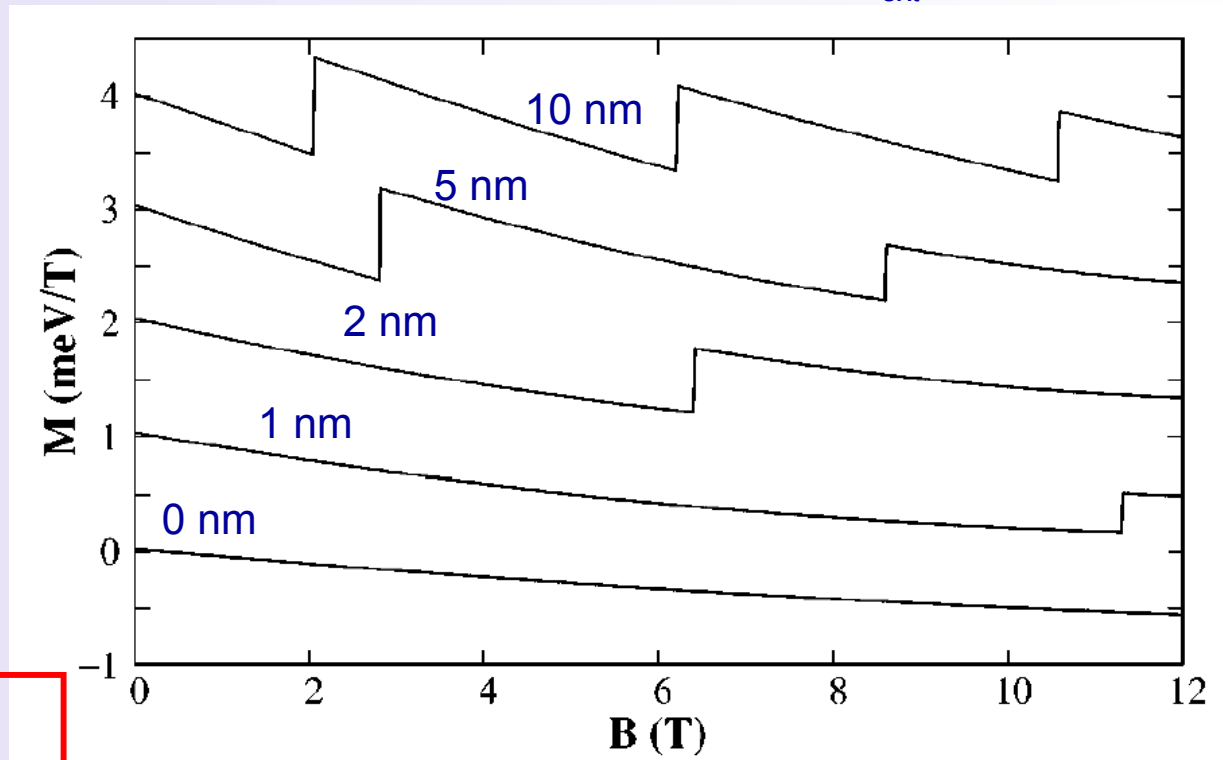
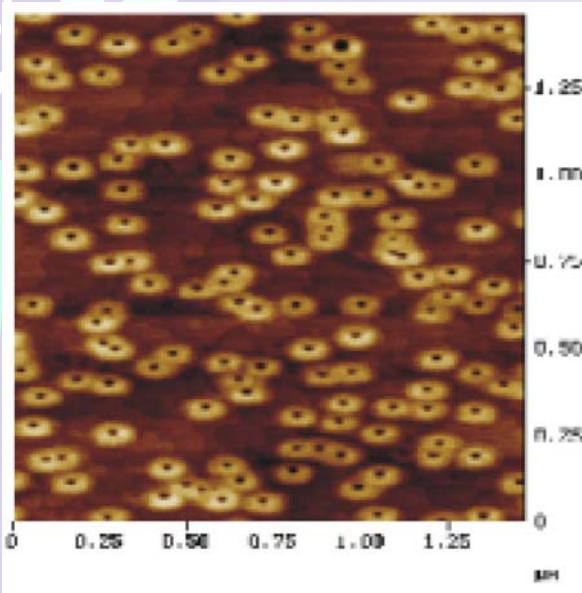
*J.M. García et al., Appl. Phys. Lett.*  
**71** (1997) 2014

*T. Raz et al., Appl. Phys. Lett* **82**  
(2003) 1706

*B.C. Lee, C.P. Lee, Nanotech.* **15**  
(2004) 848

# Magnetization: finger print

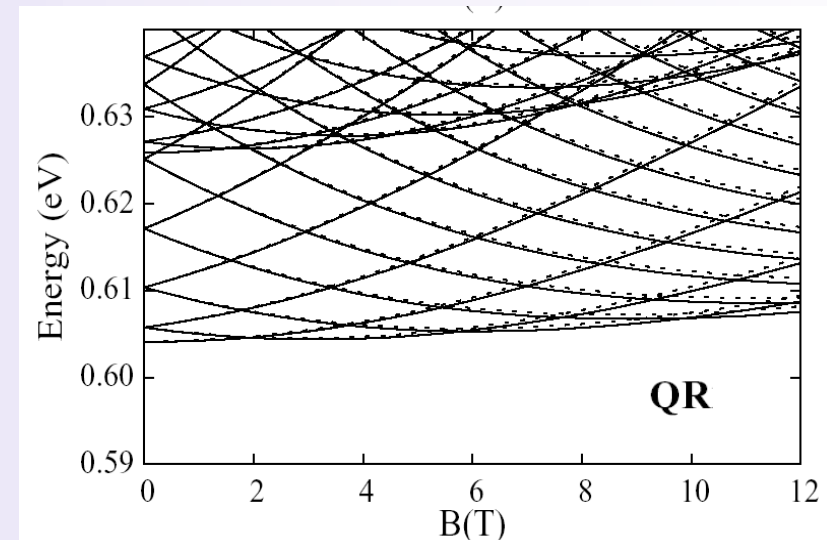
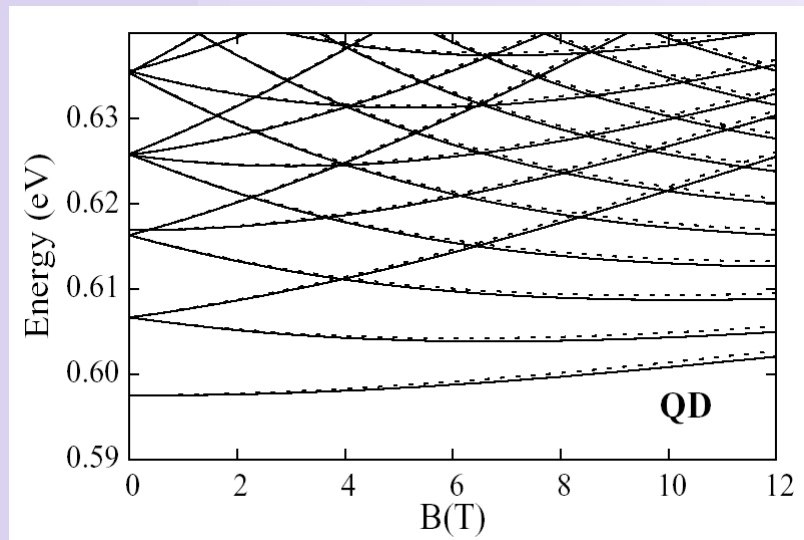
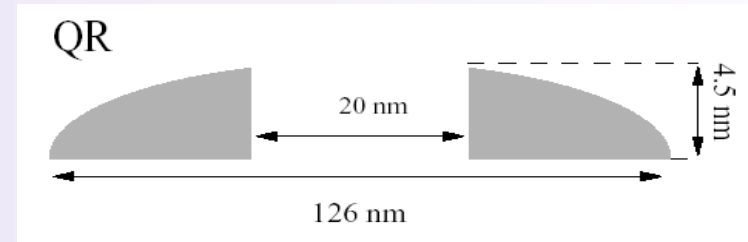
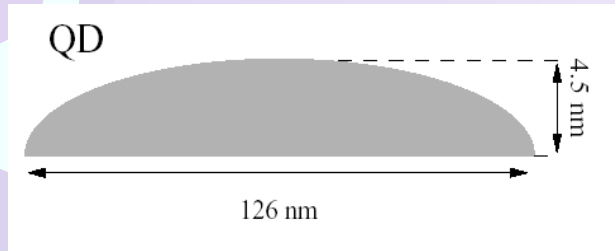
$R_{\text{ext}} = 60 \text{ nm}$



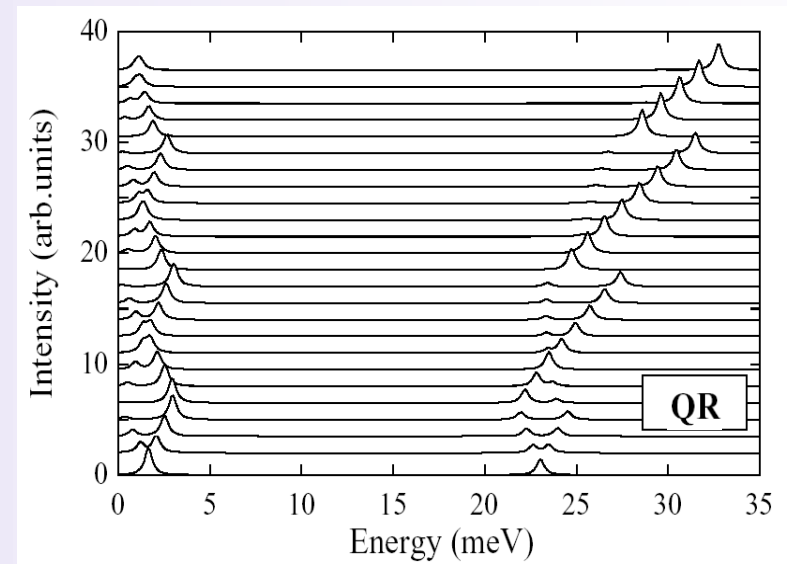
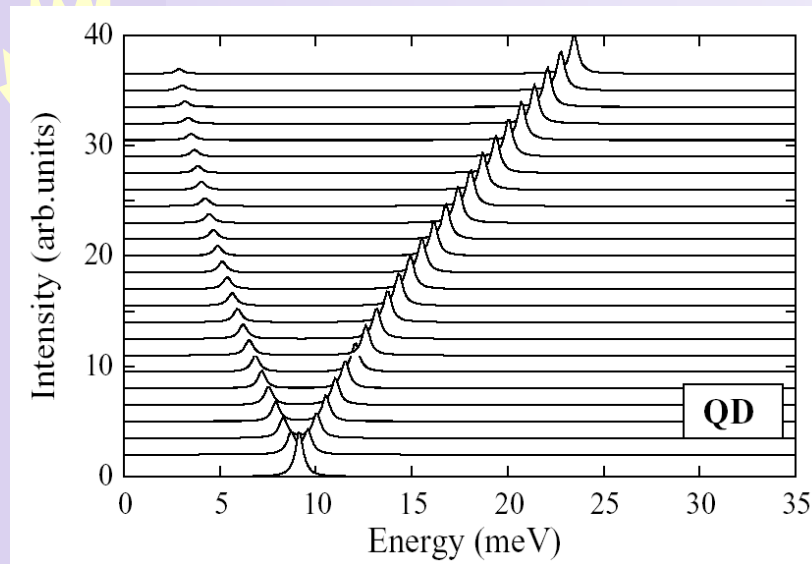
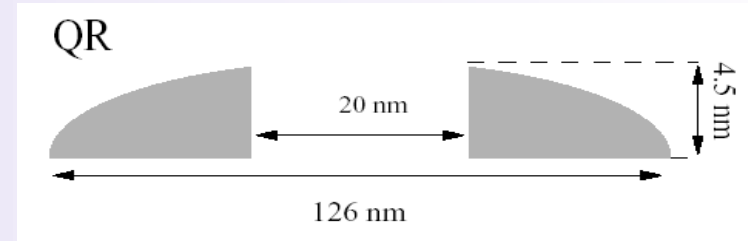
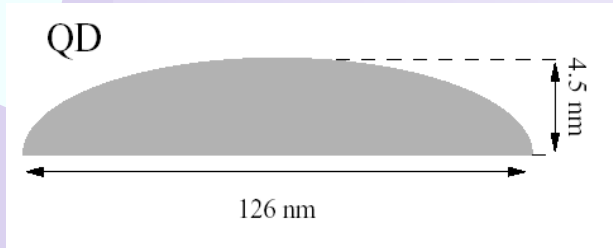
One electron magnetization ( $T=0 \text{ K}$ )

Is the ring-like geometry preserved after the sample is covered with matrix material?

# Energy vs. magnetic field of one electron in a QD and a QR



# FIR absorption of one electron in QD and QR



# Electron in a magnetic field

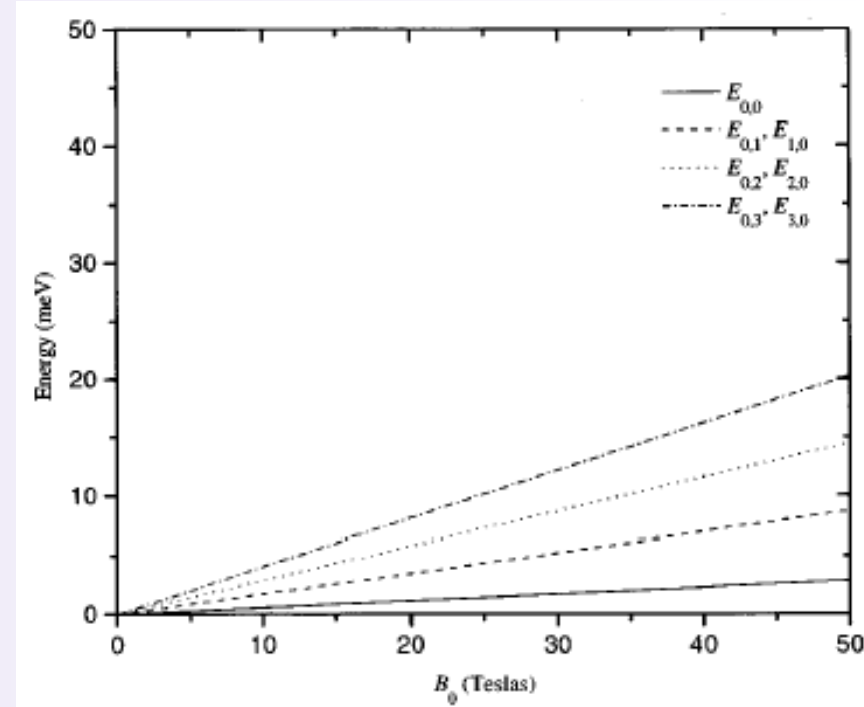
$$\hat{\mathcal{H}} = \frac{(\hat{p} - eA)^2}{2m_e} + V$$

$$\vec{B} = B_0 \vec{k}$$

$$\vec{A} = \left(-\frac{1}{2}y B_0, \frac{1}{2}x B_0, 0\right)$$

$$\begin{aligned}\hat{\mathcal{H}} &= -\frac{\hbar^2}{2m_e} \nabla^2 - \frac{eB}{2m_e} \hat{L}_z + \frac{e^2 B^2}{8m_e} \rho^2 + V \\ &= \frac{\hat{p}_z^2}{2m_e} + \hat{\mathcal{H}}_{HO}^{2D} - \frac{eB}{2m_e} \hat{L}_z + V\end{aligned}$$

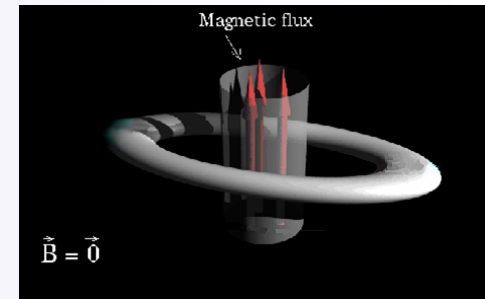
- Landau levels
- No crossings



# Electron in a potential vector but no magnetic field

$B = B_0$  if  $0 < \rho < a$  and  $B = 0$  otherwise

$$\vec{A} = A_\phi \vec{u}_\phi = \begin{cases} \frac{1}{2} B \rho \vec{u}_\phi & 0 < \rho < a \\ \frac{Ba^2}{2\rho} \vec{u}_\phi & a < \rho < \infty \end{cases}$$



$$\hat{\mathcal{H}} = \frac{(\hat{p} - eA)^2}{2m_e} + V$$

$$\hat{\mathcal{H}} = -\frac{\hbar^2}{2m_e} \nabla^2 + \frac{i\hbar e B a^2}{m_e 2\rho^2} \frac{\partial}{\partial \phi} + \frac{e^2 B^2 a^4}{8m_e \rho^2} + V$$

$$\frac{\partial}{\partial \phi} \rightarrow \frac{\partial}{\partial \phi} - \frac{i e B a^2}{2\hbar} = \frac{\partial}{\partial \phi} + i \frac{\Phi}{\Phi_0}$$

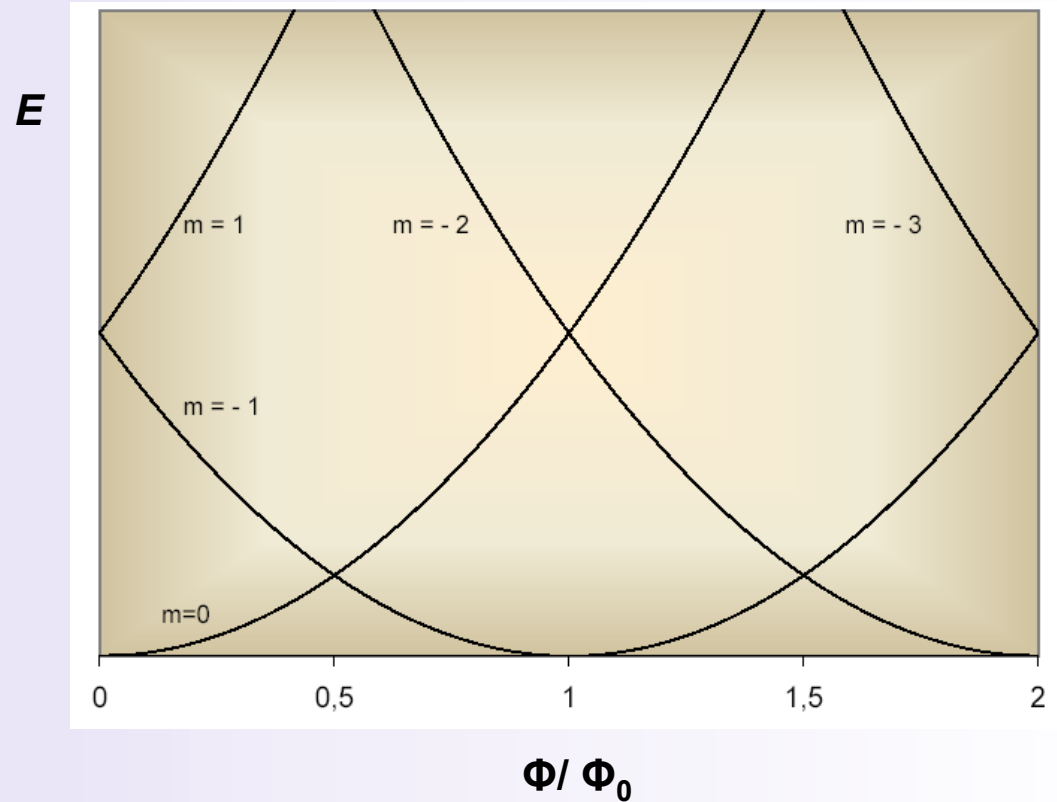


# Aharonov-Bohm Effect

1D QR

$$\hat{\mathcal{H}} = -\frac{\hbar^2}{2m_e R^2} \left( \frac{\partial}{\partial \phi} + i \frac{\Phi}{\Phi_0} \right)^2$$

$$E_m = \frac{1}{2} (m + F)^2$$



- Periodic symmetry changes of the energy levels
- Energetic oscillations

# Fractional Aharonov-Bohm Effect

$$\hat{H}(1, 2) = -\frac{1}{2R^2} \left( \frac{\partial}{\partial \phi_1} + iF \right)^2 - \frac{1}{2R^2} \left( \frac{\partial}{\partial \phi_2} + iF \right)^2 + \frac{1}{r_{12}}$$

$$r_{12} = 2R \left| \sin \frac{\phi_2 - \phi_1}{2} \right|$$

Disregarding the Coulomb interaction:

$$E(m_1, m_2) = \frac{1}{2R^2} \left[ (m_1 + F)^2 + (m_2 + F)^2 \right] \rightarrow$$

periodic changes of the ground state at the same values of flux as in the one electron case.

Eigenfunctions: singlets (S) and triplets (T)

$$|m_1, m_2; S/T\rangle = e^{im_1\phi_1} e^{im_2\phi_2} \pm e^{im_1\phi_2} e^{im_2\phi_1}$$

$|m, m; T\rangle$  does not exist (is zero)  $\rightarrow$  ground state always singlet.

# Fractional Aharonov-Bohm Effect

$$s = \frac{1}{2}(\phi_1 + \phi_2)$$

$$r = \frac{1}{2}(\phi_1 - \phi_2)$$

$$|M, m; S\rangle = e^{iMs} \cos mr$$

$$|M, m; T\rangle = e^{iMs} \sin mr$$

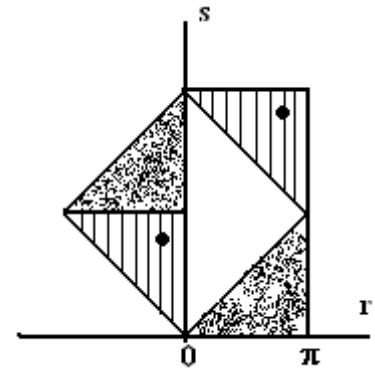
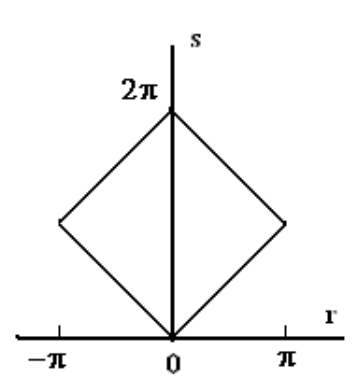
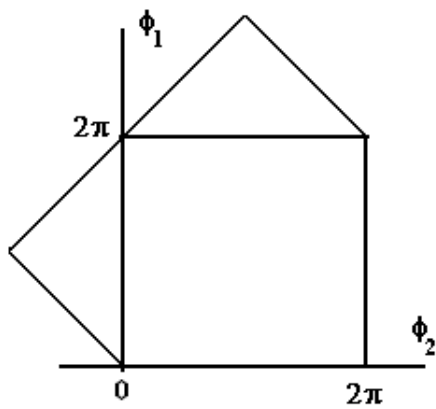
no triplets for  $m=0$

$$M = m_1 + m_2, m = m_1 - m_2$$

$M, m$  same parity

$$E(M, m) = \frac{1}{4R^2} [(M + 2F)^2 + m^2]$$

Mapping between  $(\phi_1, \phi_2)$  and  $(r, s)$  domains.



$$\phi_1 \equiv \phi_1 + 2\pi$$



$$(s, r) \equiv (s + \pi, r + \pi)$$

# Fractional Aharonov-Bohm Effect

$$\hat{\mathcal{H}}(1, 2) = -\frac{1}{4R^2} \left( \frac{\partial}{\partial s} + 2iF \right)^2 - \frac{1}{4R^2} \frac{\partial^2}{\partial r^2} = \hat{\mathcal{H}}_s + \hat{\mathcal{H}}_r$$

domain

$$0 < s < 2\pi$$

$$0 < r < \pi$$

With Coulomb interaction:

$$\hat{\mathcal{H}}_r = -\frac{1}{4R^2} \frac{\partial^2}{\partial r^2} + \frac{1}{2R|\sin r|}$$

Eigenfunctions:  $e^{iMs} \Psi_n(r)$

Pauli's principle restrictions in the presence of Coulomb interactions:

$$(s, r) \equiv (s + \pi, r + \pi)$$

$$\mathcal{P}_{12}s = s$$

$$\mathcal{P}_{12}r = -r$$

$$e^{iMs} \Psi_n(-r) = e^{iMs} e^{iM\pi} \Psi_n(\pi - r)$$



$$\hat{\mathcal{P}}_{12} \Psi_n(r) = (-1)^M \Psi_n(\pi - r)$$

$$\hat{\mathcal{P}}_{12} \Psi_n(r) = (-1)^{(M+n)} \Psi_n(r)$$



# Fractional Aharonov-Bohm Effect

$| (M, n) S/T \rangle$

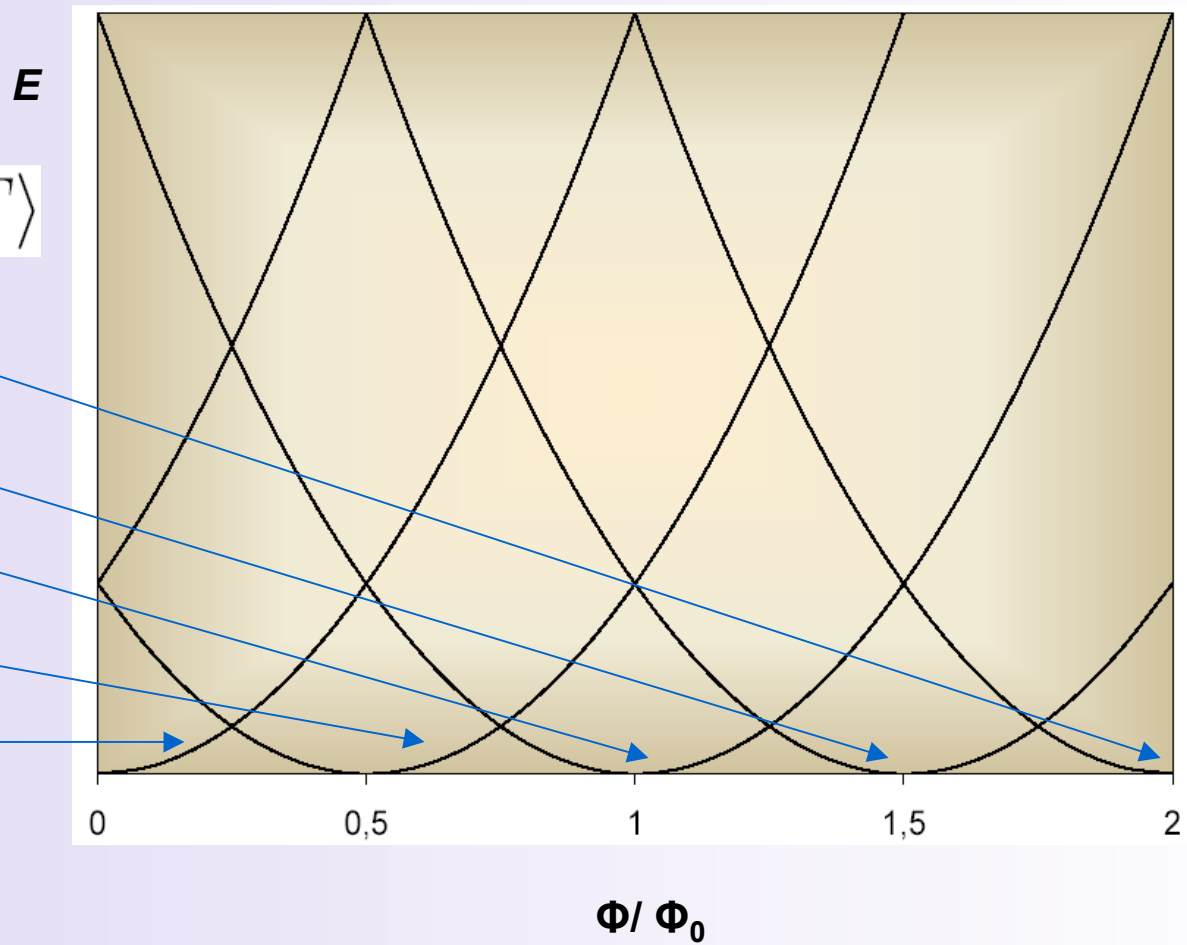
$| (-4, 0) S \rangle$

$| (-3, 0) T \rangle$

$| (-2, 0) S \rangle$

$| (-1, 0) T \rangle$

$| (0, 0) S \rangle$



# Fractional Aharonov-Bohm Effect

Coulomb interaction in a 1D system is unrealistically large.

$$\mathcal{H}_\xi = -\frac{1}{4R^2} \frac{\partial^2}{\partial r^2} + \frac{1}{\xi + 2R |\sin r|}$$

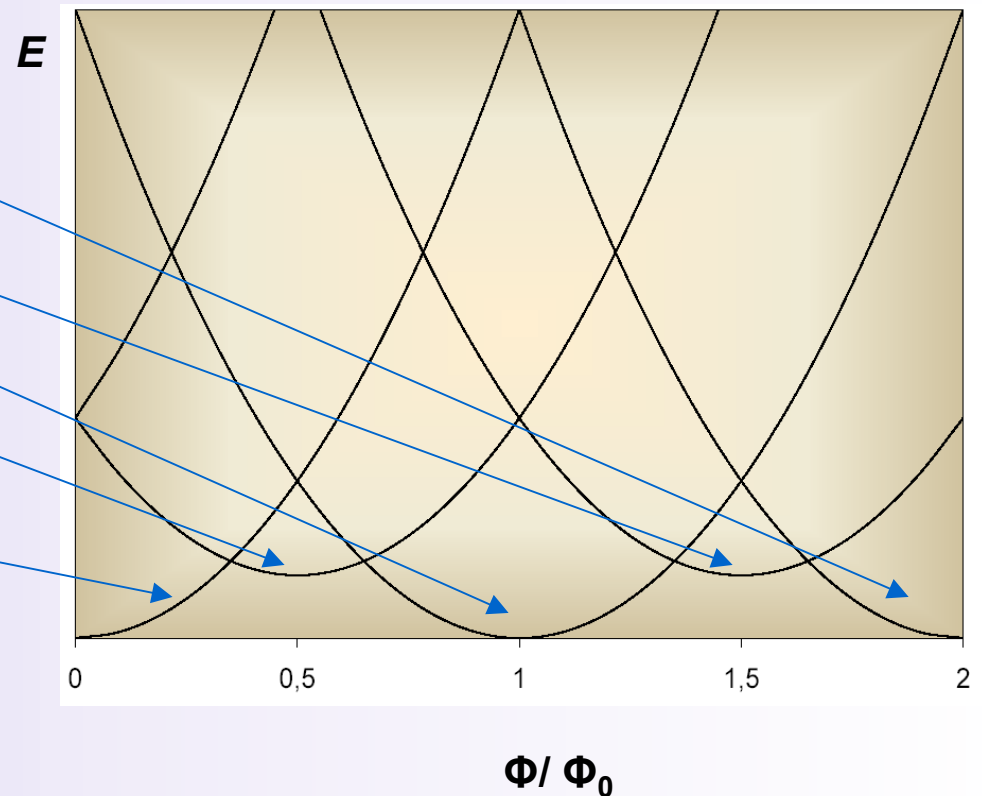
$|(-4, 0) S\rangle$

$|(-3, 0) T\rangle$

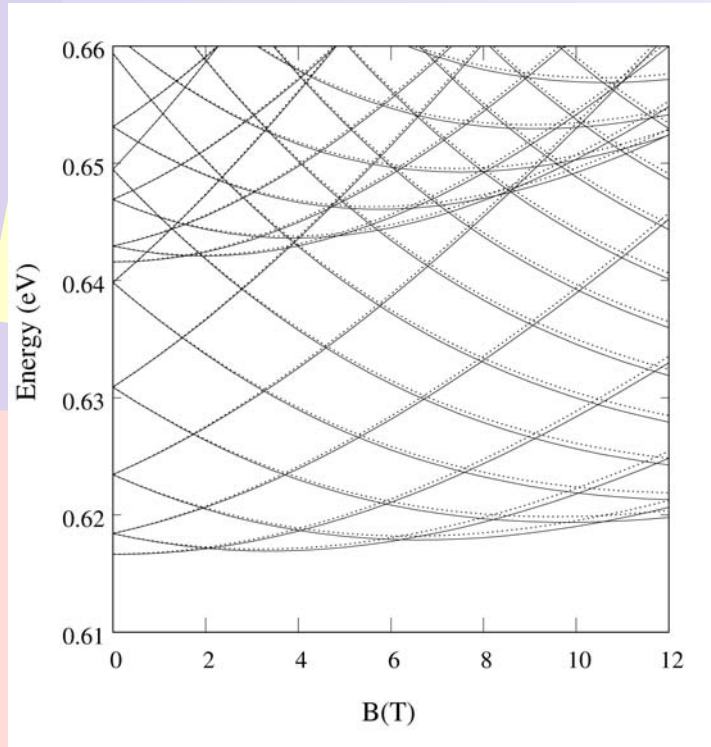
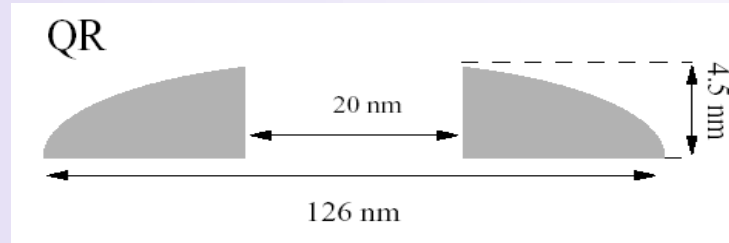
$|(-2, 0) S\rangle$

$|(-1, 0) T\rangle$

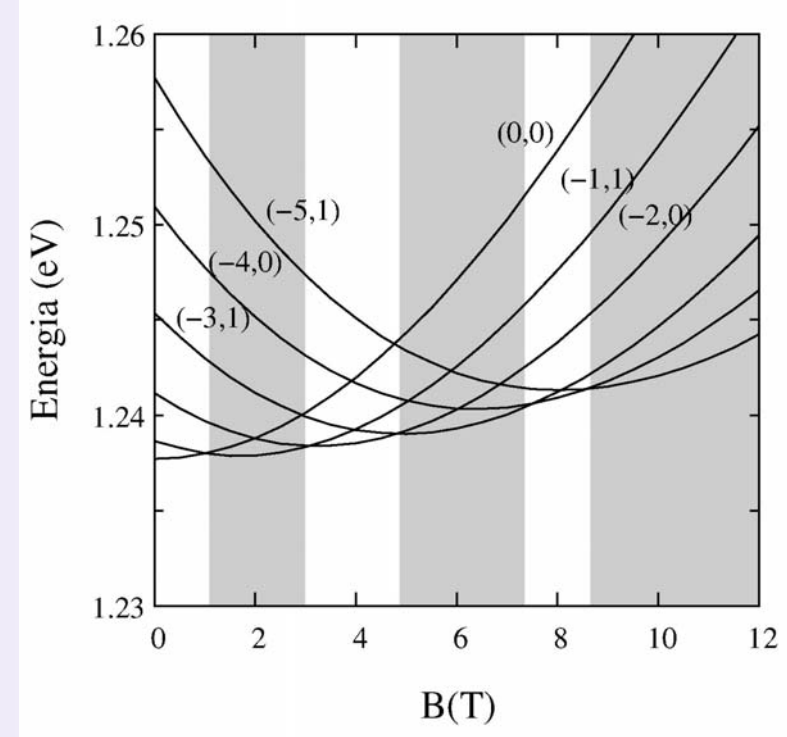
$|(0, 0) S\rangle$



# Fractional Aharonov-Bohm Effect

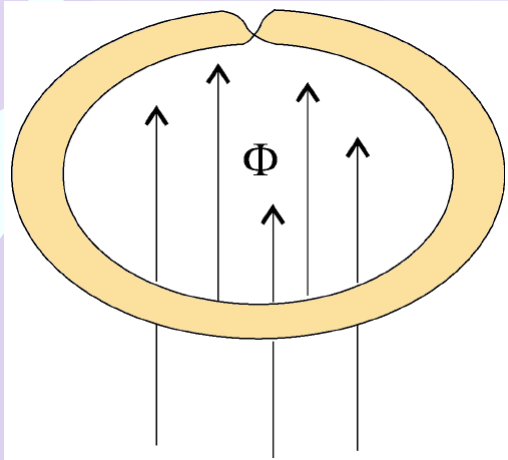


1 electron



2 electrons coulomb  
interaction

# Aharonov-Bohm Effect in a Möbius strip



Microscopic Semiconductor NbSe<sub>3</sub> Moebius strip  
S. Tanda et al., Nature **417** (2002) 397.

Möbius strip cannot be pressed into a 1D structure

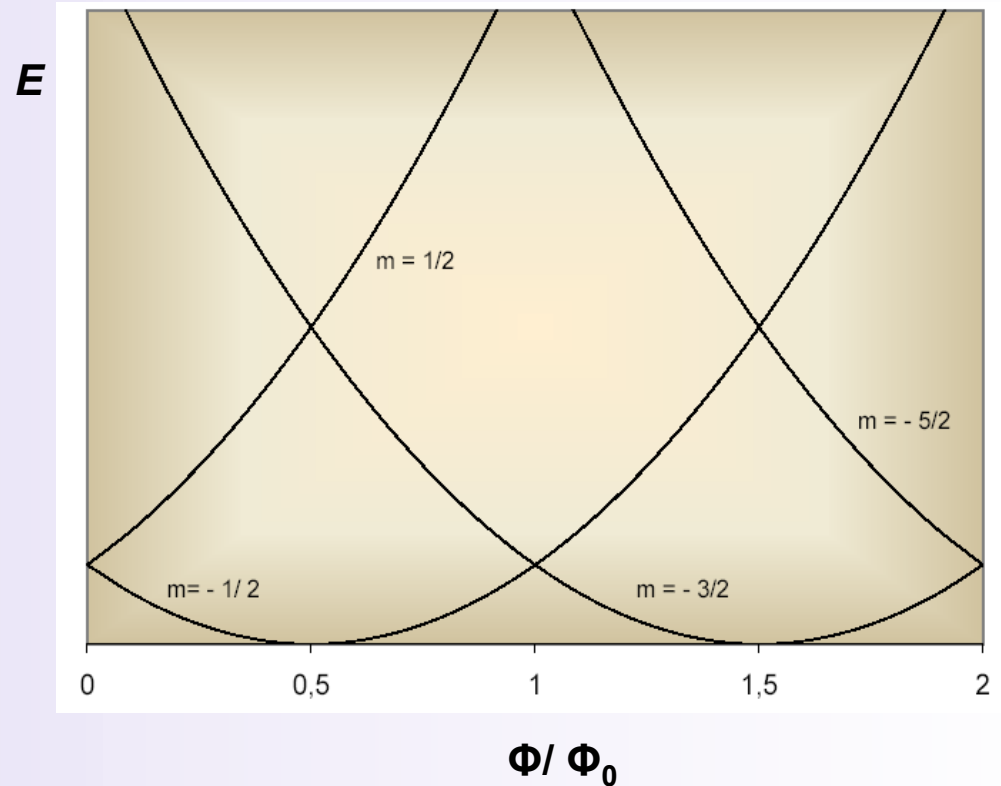
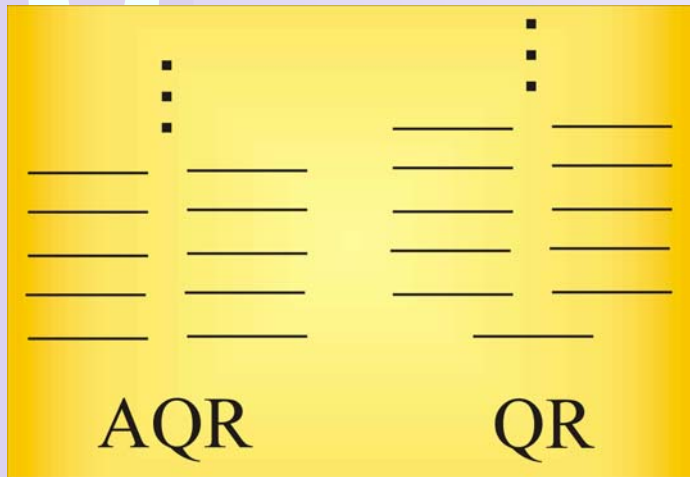
$$\rightarrow \text{1D QR } \hat{\mathcal{H}} = -\frac{\hbar^2}{2m_e R^2} \left( \frac{\partial}{\partial \phi} + i \frac{\Phi}{\Phi_0} \right)^2 \quad \rightarrow \quad E_m = \frac{1}{2} (m + F)^2$$

with antiperiodic BCs  $\Psi_m(\phi + 2\pi) = -\Psi_m(\phi)$

$$\rightarrow \quad m = \pm 1/2, \pm 3/2, \pm 5/2, \dots$$



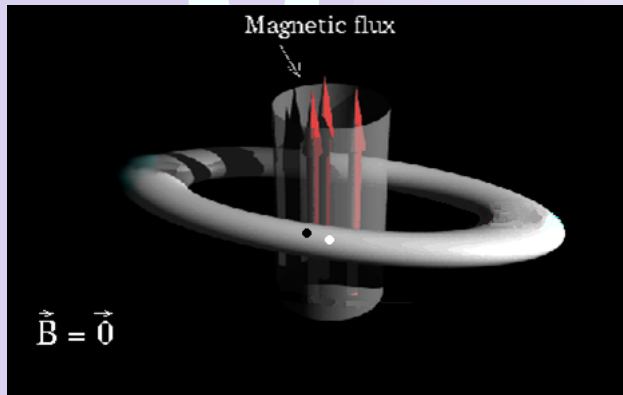
# Aharonov-Bohm Effect in a Möbius strip



$$\hat{\mathcal{P}}_{12} \Psi_n(r) = (-1)^{(M+n+1)} \Psi_n(r)$$

AQR same picture as in QR except for a shift of half flux unit.

# Optic Aharonov-Bohm effect: excitons



**Exciton is a neutral entity... Should it be sensitive to the applied magnetic flux?**

$$\hat{\mathcal{H}} = -\frac{1}{2m_e^*R^2} \left( \frac{\partial}{\partial \phi_e} + iF \right)^2 - \frac{1}{2m_h^*R^2} \left( \frac{\partial}{\partial \phi_h} - iF \right)^2 - \frac{1}{2R|\sin \frac{\phi_e - \phi_h}{2}|}$$

**Without Coulomb term:**

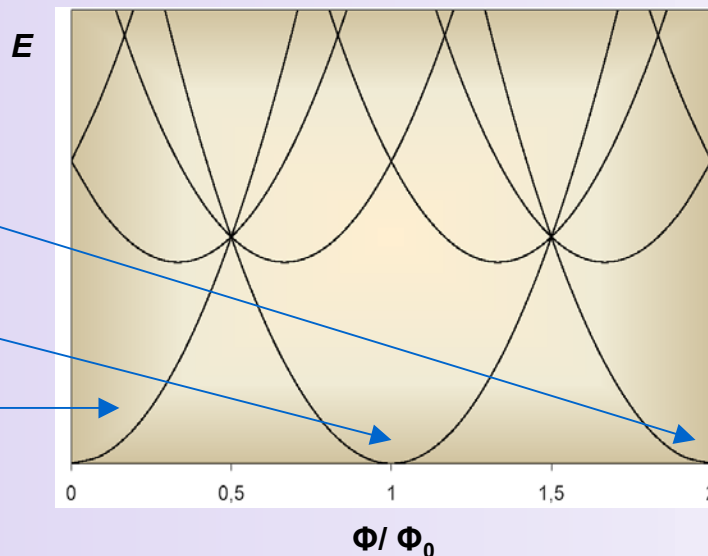
$$E - E_g = \frac{1}{2m_e^*R^2} (M_e + F)^2 + \frac{1}{2m_h^*R^2} (M_h - F)^2$$

$|M_e, M_h\rangle$

$|-2, 2\rangle$

$|-1, 1\rangle$

$|0, 0\rangle$

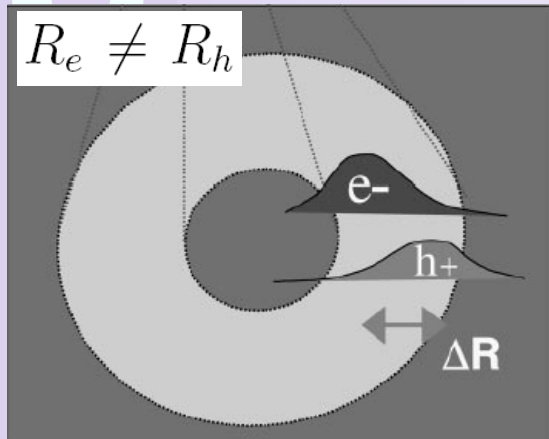


$$M_L = M_e + M_h = 0$$



**Exciton always Bright**

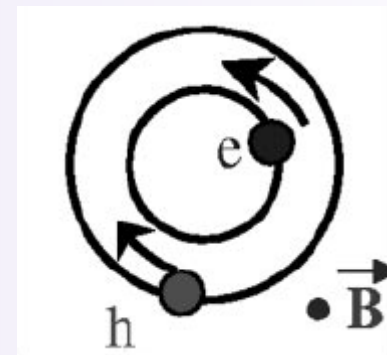
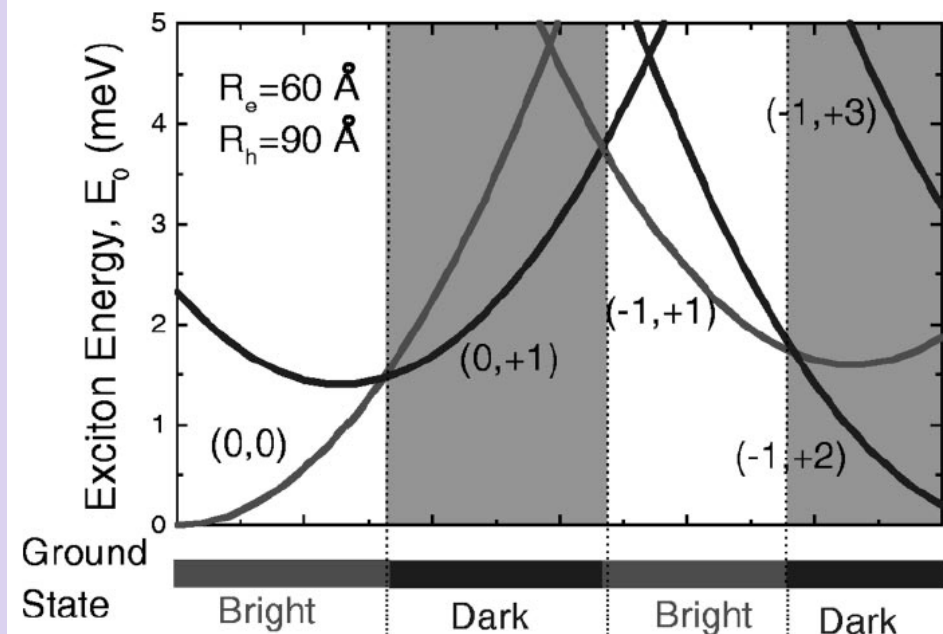
# Optic Aharonov-Bohm effect



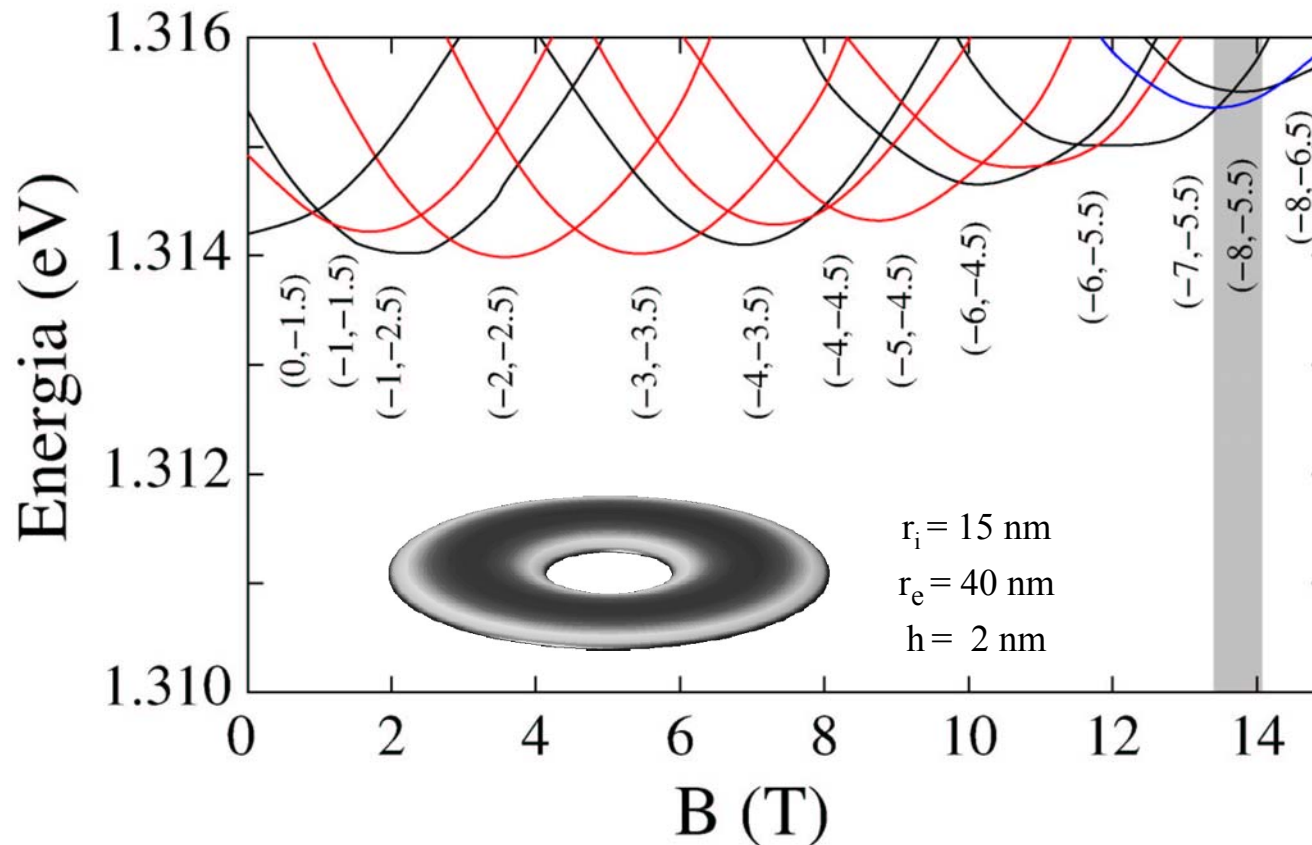
A.O. Govorov et al. *Phys. Rev. B* **68** (2003) 075307

**Without Coulomb term:**

$$E = E_g + \frac{1}{2m_e^* R_e^2} (M_e + F_e)^2 + \frac{1}{2m_h^* R_h^2} (M_h - F_h)^2$$



# Optic Aharonov-Bohm effect

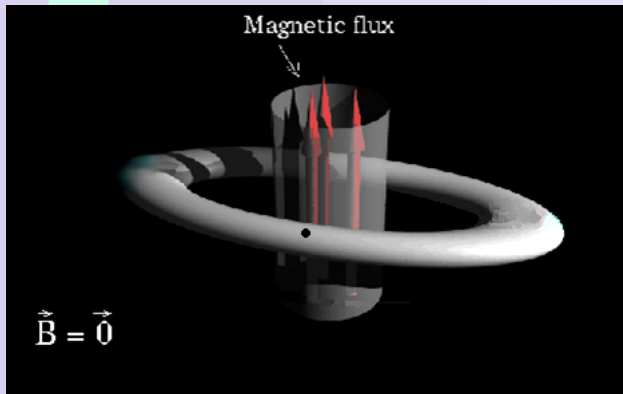


J.I. Climente, J. Planelles and W. Jaskólski, *Phys. Rev. B* **68** (2003) 075307

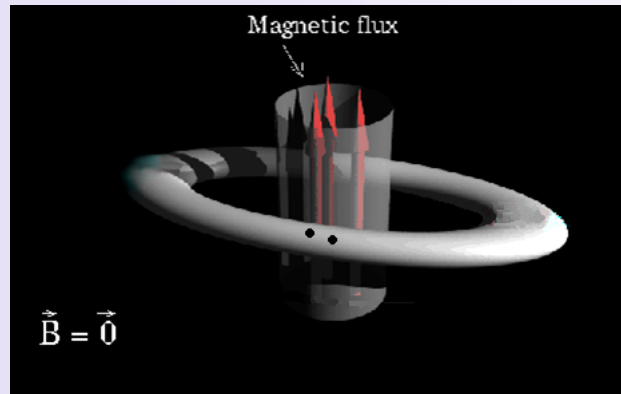
Experimentally no trace of this effect was detected for magnetic-field values up to 9 T (D. Haft et. al, *Physica E* **13**, (2002 )165)

# Aharonov-Bohm Effect

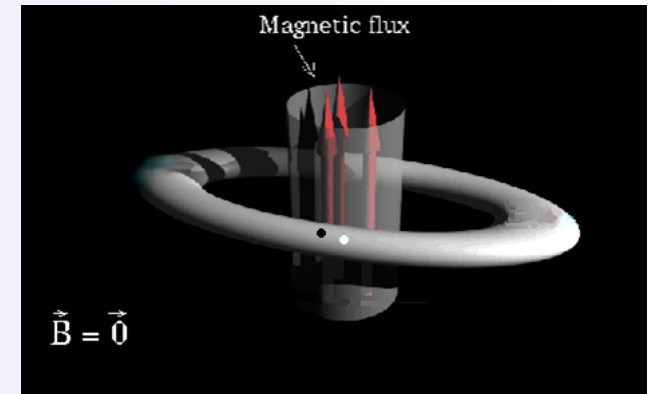
Observable pure quantum mechanical effect: Periodic oscillations



AB effect  
period  $1F$



Fractional AB  
period  $\frac{1}{2} F$



Optic AB  
Only if  $\Delta F_{e,h} \neq 0$